INTERIOR HÖLDER REGULARITY OF THE LINEARIZED MONGE-AMPÈRE EQUATION

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ABSTRACT. In this paper, we investigate the interior Hölder regularity of solutions to the linearized Monge-Ampère equation. In particular, we focus on the cases with singular right-hand side, which arise from the study of the semigeostrophic equation and singular Abreu equations. In the two-dimensional case, we give a new proof of the Caffarelli-Gutiérrez Hölder estimate (Amer. J. Math. 119 (1997), no. 2, 423-465) and the result of Le (Comm. Math. Phys. 360 (2018), no. 1, 271-305) for the linearized Monge-Ampère equation with singular right-hand side term in divergence form. The main new ingredient in the proof contains the application of the partial Legendre transform to the linearized Monge-Ampère equation. Building on this idea, we also establish a new Moser-Trudinger type inequality in dimension two. In higher dimensions, we derive the interior Hölder estimate under certain integrability assumptions on the coefficients using De Giorgi's iteration.

1. Introduction

In this paper, we investigate the interior Hölder regularity for solutions to the inhomogeneous linearized Monge-Ampère equation

(1.1)
$$\sum_{i,j=1}^{n} D_j \left(\Phi^{ij} D_i u \right) = \operatorname{div} F + f$$

in a bounded convex domain $\Omega \subset \mathbb{R}^n$ $(n \geq 2)$, where $\Phi = (\Phi^{ij})$ is the cofactor matrix of the Hessian matrix of a convex function $\phi \in C^2(\Omega)$, $F : \Omega \to \mathbb{R}^n$ is a vector field,

and $f:\Omega\to\mathbb{R}$ is a function. Since Φ is divergence free, i.e. $\sum_{j=1}^n D_j\Phi^{ij}=0$, for all

 $i = 1, 2, \dots, n$, we can rewrite (1.1) in the non-divergence form as follows:

(1.2)
$$\sum_{i,j=1}^{n} \Phi^{ij} D_{ij} u = \operatorname{div} F + f.$$

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When the equation is uniformly elliptic, it is well known that the Harnack and Hölder estimates of the solution are established in the classical De Giorgi-Nash-Moser theory for equations of divergence form [De, Na, Mo1], and in the Krylov-Safonov theory for general equations of non-divergence form [KS]. The main interest on equation (1.1) or (1.2) lies in the lack of uniform ellipticity. In a celebrated work, Caffarelli-Gutiérrez obtained the Harnack inequality and the Hölder estimate for equation (1.1) or (1.2) with $F \equiv 0$ and $f \equiv 0$, under the \mathcal{A}_{∞} condition [CG]. In particular, this condition is satisfied if

(1.3)
$$0 < \lambda \le \det D^2 \phi \le \Lambda \quad \text{in } \Omega.$$

The Hölder estimate for the inhomogeneous equation $(F = 0 \text{ and } f \neq 0)$, as well as higher order estimates and the boundary regularity, were later established by [TrW, GN1, GN2, LN1, LS] under certain assumptions on f. For further extensions and related work, one can refer to [Le1, Le5, LN2, KLWZ] and the references therein.

When $F \neq 0$, the equation arises from the study of semigeostrophic equations [ACDF, Le2, Lo] and singular Abreu equations [KLWZ, Le4, LZ] in the study of convex functionals with a convexity constraint related to the Rochet-Choné model for the monopolist problem in economics. So far, very little is known about the regularity of (1.1) when $F \neq 0$. We focus mainly on the linearized Monge-Ampère equation under the condition (1.3). Loeper [Lo] obtained the interior Hölder regularity for (1.1) under the stronger assumption that $\det D^2 \phi$ is sufficiently close to a positive constant, using the $W^{2,p}$ estimate of the Monge-Ampère equation and a result derived in [Tr2]. Later, Le [Le2] showed the same result when n=2, only assuming (1.3). The main ingredient used in [Le2] is the $W^{2,1+\varepsilon}$ -estimate of the Monge-Ampère equation established by [DFS, Sc]. Under certain integral bounds on the Hessian $D^2 \phi$, Le [Le6] also extended the Hölder estimates to higher dimensions. More recently, Kim [Kim] derived similar estimates for (1.1) with drift terms under a similar condition as [Le6]. For the estimates to boundary case, see [Le3].

The results in [Kim, Le2, Le6, Lo] used both the De Giorgi-Nash-Moser iteration and the Caffarelli-Gutiérrez estimate, corresponding to the divergence form and the non-divergence form of linearized Monge-Ampère equations, respectively. Note that it has been pointed out that in general it is impossible to obtain the Caffarelli-Gutiérrez estimate by the De Giorgi-Nash-Moser iteration [TW, Remark 3.4], which means that the celebrated theory of Caffarelli-Gutiérrez's is essential in their arguments. However, in this paper we find that there is a new proof of the theorem below without using the Caffarelli-Gutiérrez estimate in dimension two.

Theorem 1.1. Assume n=2. Let $\phi \in C^2(\Omega)$ be a convex function satisfying (1.3). Let $F:=(F^1(x),F^2(x)):\Omega \to \mathbb{R}^2$ be a bounded vector field and $f\in L^r(\Omega)$ for r>1. Given $\Omega'\subset\subset\Omega$ and $p\in(0,+\infty)$, then for every solution u to (1.1) in Ω , there is

$$||u||_{C^{\gamma}(\Omega')} \le C \left(||u||_{L^{p}(\Omega)} + ||F||_{L^{\infty}(\Omega)} + ||f||_{L^{r}(\Omega)} \right),$$

where constant $\gamma > 0$ depending only on λ and Λ , and constant C > 0 depending only on p, r, λ , Λ , and dist(Ω' , $\partial\Omega$).

Note that although we assume $\phi \in C^2(\Omega)$, the derived estimates are independent of the smoothness of ϕ and depend only on the structure constants. Theorem 1.1 includes [Le2, Theorem 1.3] and Caffarelli-Gutiérrez's estimate [CG] in dimension two. Our main new idea in Theorem 1.1 is the use of the partial Legendre transform. After the partial Legendre transform, (1.1) becomes a linear uniformly elliptic equation in divergence form with singular right-hand side (see (2.7)). Thus, the De Giorgi-Nash-Moser theory implies that the solution after transformation is Hölder continuous. Then transforming back to the original solution gives us the result. We still need the $W^{2,1+\varepsilon}$ -estimate of the Monge-Ampère equation to guarantee that the condition in De Giorgi-Nash-Moser's theory is satisfied. The partial Legendre transform has been widely used in the study of the Monge-Ampère equation [DS, Fi, GP, Liu], and it has also been used recently to study the Monge-Ampère type fourth order equation [LZ, WZ]. However, we didn't find its use in the linearized Monge-Ampère equation. Our proof can be seen as an attempt in this direction.

On the other hand, due to the divergence form of the equation (1.1), we already know that some interesting Sobolev inequalities of Monge-Ampère type were obtained by [TW] in dimension $n \geq 3$ and [Le2] in dimension n = 2 (see [Ma1] for some extensions and [WZ23] for a complex version). Since we will use the Monge-Ampère type Sobolev inequality later, we restate it here.

Theorem 1.2 ([TW, Theorem 1.1], [Le2, Proposition 2.6]). Assume $n \geq 2$. Let $\phi \in C^2(\Omega)$ be a convex function satisfying (1.3). Then there exists a constant $C_{Sob} > 0$, depending only on n, λ , Λ , and Ω such that

$$\left(\int_{\Omega} |u|^{2^*} dx\right)^{\frac{1}{2^*}} \le C_{Sob} \left(\int_{\Omega} \Phi^{ij} D_i u D_j u dx\right)^{\frac{1}{2}}, \quad \forall u \in C_0^{\infty}(\Omega),$$

where $2^* = \frac{2n}{n-2}$ for $n \ge 3$, and any $2^* > 2$ for n = 2.

Therefore, with the new idea in Theorem 1.1, we also establish a new Moser-Trudinger type inequality in two dimensions. To simplify the notation, we write

$$||Du||_{\Phi}^2 := \int_{\Omega} \Phi^{ij} D_i u D_j u \, \mathrm{d}x.$$

Theorem 1.3. Let Ω be a uniformly convex domain in \mathbb{R}^2 and $\phi \in C^2(\Omega)$ be a convex function satisfying (1.3). Assume that $\phi|_{\partial\Omega}$ and $\partial\Omega$ are of class C^3 . For any $u \in C_0^{\infty}(\Omega)$, there exists a constant C > 0 depending only on λ , $\Lambda \|\phi\|_{C^3(\partial\Omega)}$, the uniform convexity radius of $\partial\Omega$ and the C^3 regularity of $\partial\Omega$ such that

(1.4)
$$\int_{\Omega} e^{\beta \frac{u^2}{\|Du\|_{\Phi}^2}} dx_1 dx_2 \le C|\Omega|^{\frac{\varepsilon_0}{2+\varepsilon_0}},$$

where $\beta \leq 4\pi \frac{1+\varepsilon_0}{2+\varepsilon_0} \min\{\lambda,1\}$, and ε_0 depending only on λ and Λ is obtained by the global $W^{2,1+\varepsilon}$ -estimate for Monge-Ampère equations.

Remark 1.4. If $\phi(x) = \frac{1}{2}|x|^2$, we know that $\{\Phi^{ij}\} = id$, $\lambda = 1$ and $\varepsilon_0 = +\infty$. So in this case the inequality (1.4) is identical to the classical Moser-Trudinger inequality.

In higher dimensions, very little is known when the singular term F appears. For general degenerate linear elliptic equations, there are some extensions of the classical De Giorgi-Nash-Moser theory with some integral conditions on the elliptic coefficients in [MS, Tr2]. Therefore, in the second part of this paper, we will investigate the interior regularity in terms of (1.1) in higher dimensions with certain assumptions. As mentioned earlier, directly applying the De Giorgi-Nash-Moser iteration faces challenges due to the lack of a suitable estimate for $|D^2\phi|$. Therefore, we need to introduce further assumptions on F to allow the use of the De Giorgi-Nash-Moser iteration. These assumptions are specified in the following theorem. Denote

$$S_{\phi}(x,h) := \{ y \in \Omega \mid \phi(y) < \phi(x) + D\phi(x) \cdot (y-x) + h \}$$

as the section of ϕ centered at $x \in \Omega$ with height h > 0.

Theorem 1.5. Let $\phi \in C^2(\Omega)$ be a convex function satisfying (1.3). For q > n, let $F := (F^1(x), \dots, F^n(x)) : \Omega \to \mathbb{R}^n$ be a vector field satisfying $||F_{\phi}||_{L^q(\Omega)} < \infty$, where $F_{\phi}(x) := (D^2\phi(x))^{1/2} F$ and $f : \Omega \to \mathbb{R}$ be a function satisfying $||f||_{L^{q_*}(\Omega)} < \infty$, where $q_* = \frac{nq}{n+q}$. Given a section $S_{\phi}(x_0, 2h_0) \subset \Omega$, for every solution u to (1.1) in $S_{\phi}(x_0, 2h_0)$ and for all $x \in S_{\phi}(x_0, h_0)$, there is

 $|u(x) - u(x_0)| \le C \left(||u||_{L^{\infty}(S_{\phi}(x_0, 2h_0))} + ||F_{\phi}||_{L^q(S_{\phi}(x_0, 2h_0))} + ||f||_{L^{q_*}(S_{\phi}(x_0, 2h_0))} \right) |x - x_0|^{\gamma},$ where constant $\gamma > 0$ depending only on n, λ and Λ , and constant C > 0 depending only on n, q, λ , Λ , h_0 and diam(Ω).

Remark 1.6. Some remarks in order.

- (1) With specific assumptions regarding the integral bound of $D^2\phi$, one can also get a result similar to Theorem 1.5 in all dimensions, see details in [Le6, Theorem 15.6] and [Kim, Corollary 1.2].
- (2) It would be interesting to remove the assumption on F_{ϕ} and/or the integral bound of $D^2\phi$ in Theorem 1.5 and instead simply assume that (1.3) and F is bounded. But so far we haven't come up with a way to deal with this when n > 3.

The framework of the proof of Theorem 1.5 is similar to that of [Le2, Lo]. First, we need to derive a new weak maximum principle (Theorem 3.1 and Corollary 3.2) for the solution of (1.1) by De Giorgi's iteration. Combining the new weak maximum principle with the Caffarelli-Gutiérrez's Harnack inequality then gives Theorem 1.5.

Remark 1.7. We can consider a more general class of linearized Monge-Ampère equations given by:

where A(x) is a symmetric matrix satisfying $\lambda \Phi \leq A \leq \Lambda \Phi$. This equation, (1.5), has a similar divergence structure to the previously studied equation (1.1). Using the methods outlined in this paper or that in [Le2, Kim], we can obtain interior Hölder estimates for (1.5) under the same assumptions. One motivation for studying (1.5)

arises from the singular Abreu equation. Through direct computations (see [KLWZ, Lemma 2.1]), it is shown that the singular Abreu equation

$$U^{ij}D_{ij}\left[(\det D^2u)^{-1}\right] = -\operatorname{div}(|Du|^{q-2}Du) + f,$$

is equivalent to

$$D_j(u^{ij}D_i\zeta) = -f,$$

where

$$\zeta = \log \det D^2 u - \frac{1}{q} |Du|^q.$$

In this case, we can set $A = (D^2u)^{-1}$ in (1.5), illustrating that (1.5) encompasses this specific case of the singular Abreu equation. This connection allows us to apply our general results on interior Hölder estimates to the singular Abreu equation, providing valuable regularity insights. This type of equation, (1.5), has been also studied by Maldonado in [Ma2].

The rest of the paper is organized as follows. In Section 2, we first apply the partial Legendre transform to (1.1), then give the proof of Theorem 1.1. The Moser-Trudinger type inequality is also proved in this section. Next, in Section 3, we present some estimates for linearized Monge-Ampère equations and then proceed to prove Theorem 1.5.

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2. Linearized Monge-Ampère equations in dimension two

In this section, we present a new proof of the interior estimate for (1.1) without Caffarelli-Gutiérrez's theory and establish a new Moser-Trudinger type inequality in dimension two.

2.1. The new equation under partial Legendre transform. In this subsection, we first derive the new equation under the partial Legendre transform. Let $\Omega \subset \mathbb{R}^2$ and $\phi(x_1,x_2)$ be a convex function on Ω . The partial Legendre transform in the x_1 -variable is

(2.1)
$$\phi^*(\xi, \eta) = \sup\{x_1 \xi - \phi(x_1, \eta)\},\$$

where the supremum is taken with respect to x_1 on the slice η is the fixed constant, namely for all x_1 such that $(x_1, \eta) \in \Omega$. This definition is taken from [Liu]. Hence, when $\phi \in C^2(\Omega)$ is a strictly convex function, we will have a injective mapping \mathcal{P} satisfying

(2.2)
$$(\xi, \eta) = \mathcal{P}(x_1, x_2) := (\phi_{x_1}, x_2) \in \mathcal{P}(\Omega) := \Omega^*,$$

where $\phi_{x_1} := D_{x_1} \phi$. In this situation, we know that

$$\phi^{\star}(\xi,\eta) = x_1 \phi_{x_1}(x_1, x_2) - \phi(x_1, x_2).$$

Indeed, it just needs ϕ to be strictly convex respect to x_1 -variable [GP]. Then a direct calculation yields

$$\frac{\partial(\xi,\eta)}{\partial(x_1,x_2)} = \begin{pmatrix} \phi_{x_1x_1} & \phi_{x_1x_2} \\ 0 & 1 \end{pmatrix}, \quad \text{and} \quad \frac{\partial(x_1,x_2)}{\partial(\xi,\eta)} = \begin{pmatrix} \frac{1}{\phi_{x_1x_1}} & -\frac{\phi_{x_1x_2}}{\phi_{x_1x_1}} \\ 0 & 1 \end{pmatrix}.$$

Hence,

$$\phi_{\varepsilon}^{\star} = x_1, \quad \phi_{\eta}^{\star} = -\phi_{x_2},$$

(2.4)
$$\phi_{\xi\xi}^{\star} = \frac{1}{\phi_{x_1x_1}}, \quad \phi_{\eta\eta}^{\star} = -\frac{\det D^2\phi}{\phi_{x_1x_1}}, \quad \phi_{\xi\eta}^{\star} = -\frac{\phi_{x_1x_2}}{\phi_{x_1x_1}}.$$

Then we know that ϕ^* is a solution to

(2.5)
$$(\det D^2 \phi) \phi_{\xi\xi}^* + \phi_{\eta\eta}^* = 0.$$

In order to derive the equation under the partial Legendre transform, we consider the associated functionals of (1.1)

(2.6)
$$A(u) := \int_{\Omega} \Phi^{ij} D_i u D_j u - 2F^i D_i u + 2f u \, \mathrm{d}x,$$

where the repeated indices are summed. Denote $\widetilde{u}(\xi,\eta) := u(\phi_{\xi}^{\star},\eta)$, $\widetilde{F}(\xi,\eta) := F(\phi_{\xi}^{\star},\eta)$ and $\widetilde{f}(\xi,\eta) := f(\phi_{\xi}^{\star},\eta)$, then we have the following equation for \widetilde{u} .

Proposition 2.1. Assume n = 2. Let u be a solution to (1.1), then \widetilde{u} satisfies

$$(2.7) \qquad \left(-\frac{\phi_{\eta\eta}^{\star}}{\phi_{\xi\xi}^{\star}}\widetilde{u}_{\xi}\right)_{\xi} + \widetilde{u}_{\eta\eta} = \left(\widetilde{F}^{1} - \widetilde{F}^{2}\phi_{\xi\eta}^{\star}\right)_{\xi} + \left(\widetilde{F}^{2}\phi_{\xi\xi}^{\star}\right)_{\eta} + \widetilde{f}\phi_{\xi\xi}^{\star} \quad in \ \Omega^{\star}.$$

Proof. Note that in dimension two, the cofactor matrix (Φ_{ij}) is

$$\begin{pmatrix} \phi_{x_2x_2} & -\phi_{x_1x_2} \\ -\phi_{x_1x_2} & \phi_{x_1x_1} \end{pmatrix},$$

then (2.6) becomes

$$A(u) = \int_{\Omega} \phi_{x_2 x_2}(u_{x_1})^2 - 2\phi_{x_1 x_2} u_{x_1} u_{x_2} + \phi_{x_1 x_1}(u_{x_2})^2 - 2F^1 u_{x_1} - 2F^2 u_{x_2} + 2fu \, dx.$$

Note that

(2.8)
$$u_{x_1} = \frac{1}{\phi_{\xi\xi}^*} \widetilde{u}_{\xi}, \ u_{x_2} = -\frac{\phi_{\xi\eta}^*}{\phi_{\xi\xi}^*} \widetilde{u}_{\xi} + \widetilde{u}_{\eta}, \ \mathrm{d}x_1 \mathrm{d}x_2 = \phi_{\xi\xi}^* \mathrm{d}\xi \mathrm{d}\eta,$$

we have

$$\begin{split} A(u) &= \int_{\Omega^{\star}} \left[-\frac{\det D^2 \phi^{\star}}{\phi_{\xi\xi}^{\star}} \left(\frac{\widetilde{u}_{\xi}}{\phi_{\xi\xi}^{\star}} \right)^2 + 2 \frac{\phi_{\xi\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \frac{\widetilde{u}_{\xi}}{\phi_{\xi\xi}^{\star}} \left(-\frac{\phi_{\xi\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \widetilde{u}_{\xi} + \widetilde{u}_{\eta} \right) \right. \\ &\quad + \frac{1}{\phi_{\xi\xi}^{\star}} \left(-\frac{\phi_{\xi\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \widetilde{u}_{\xi} + \widetilde{u}_{\eta} \right)^2 - 2\widetilde{F}^1 \frac{\widetilde{u}_{\xi}}{\phi_{\xi\xi}^{\star}} - 2\widetilde{F}^2 \left(-\frac{\phi_{\xi\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \widetilde{u}_{\xi} + \widetilde{u}_{\eta} \right) + 2\widetilde{f}\widetilde{u} \right] \phi_{\xi\xi}^{\star} \, \mathrm{d}\xi \mathrm{d}\eta \\ &= \int_{\Omega^{\star}} \left[\left(-\frac{\det D^2 \phi^{\star}}{\phi_{\xi\xi}^{\star}^2} - 2\frac{\phi_{\xi\eta}^{\star}}{\phi_{\xi\xi}^{\star}^2} + \frac{\phi_{\xi\eta}^{\star}}{\phi_{\xi\xi}^{\star}^2} \right) \widetilde{u}_{\xi}^2 + 2\frac{\phi_{\xi\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \widetilde{u}_{\xi} \widetilde{u}_{\eta} - 2\frac{\phi_{\xi\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \widetilde{u}_{\xi} \widetilde{u}_{\eta} + \widetilde{u}_{\eta}^2 \right. \\ &\quad - 2\widetilde{F}^1 \widetilde{u}_{\xi} + 2\widetilde{F}^2 \phi_{\xi\eta}^{\star} \widetilde{u}_{\xi} - 2\widetilde{F}^2 \phi_{\xi\xi}^{\star} \widetilde{u}_{\eta} + 2\widetilde{f} \phi_{\xi\xi}^{\star} \widetilde{u} \right] \mathrm{d}\xi \mathrm{d}\eta \\ &= \int_{\Omega^{\star}} \left(-\frac{\phi_{\eta\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \widetilde{u}_{\xi}^2 + \widetilde{u}_{\eta}^2 - 2\widetilde{F}^1 \widetilde{u}_{\xi} + 2\widetilde{F}^2 \phi_{\xi\eta}^{\star} \widetilde{u}_{\xi} - 2\widetilde{F}^2 \phi_{\xi\xi}^{\star} \widetilde{u}_{\eta} + 2\widetilde{f} \phi_{\xi\xi}^{\star} \widetilde{u} \right) \mathrm{d}\xi \mathrm{d}\eta =: A^{\star}(\widetilde{u}). \end{split}$$

Since u is a critical point of the functional A(u), we know that \tilde{u} is a critical point of the functional $A^*(\widetilde{u})$. Thus, it suffices to derive the Euler-Lagrange equation of $A^{\star}(\widetilde{u})$. See [WZ] for the similar argument for the Monge-Ampère type fourth order equation.

For $\varphi \in C_0^{\infty}(\Omega^*)$, by integration by parts, we have

$$\begin{split} & \frac{\mathrm{d}A^{\star}(\widetilde{u}+t\varphi)}{\mathrm{d}t} \bigg|_{t=0} \\ &= 2 \int_{\Omega^{\star}} \left(-\frac{\phi_{\eta\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \widetilde{u}_{\xi} \varphi_{\xi} + \widetilde{u}_{\eta} \varphi_{\eta} - \widetilde{F}^{1} \varphi_{\xi} + \widetilde{F}^{2} \phi_{\xi\eta}^{\star} \varphi_{\xi} - \widetilde{F}^{2} \phi_{\xi\xi}^{\star} \varphi_{\eta} + \widetilde{f} \phi_{\xi\xi}^{\star} \varphi \right) \mathrm{d}\xi \mathrm{d}\eta \\ &= 2 \int_{\Omega^{\star}} \left[-\left(-\frac{\phi_{\eta\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \widetilde{u}_{\xi} \right)_{\xi} - \widetilde{u}_{\eta\eta} + \left(\widetilde{F}^{1} \right)_{\xi} - \left(\widetilde{F}^{2} \phi_{\xi\eta}^{\star} \right)_{\xi} + \left(\widetilde{F}^{2} \phi_{\xi\xi}^{\star} \right)_{\eta} + \widetilde{f} \phi_{\xi\xi}^{\star} \right] \varphi \, \mathrm{d}\xi \mathrm{d}\eta. \end{split}$$

Then we know that the Euler-Lagrange equation of $A^*(\widetilde{u})$ is

$$-\left(-\frac{\phi_{\eta\eta}^{\star}}{\phi_{\xi\xi}^{\star}}\widetilde{u}_{\xi}\right)_{\xi}-\widetilde{u}_{\eta\eta}+\left(\widetilde{F}^{1}\right)_{\xi}-\left(\widetilde{F}^{2}\phi_{\xi\eta}^{\star}\right)_{\xi}+\left(\widetilde{F}^{2}\phi_{\xi\xi}^{\star}\right)_{\eta}+\widetilde{f}\phi_{\xi\xi}^{\star}=0,$$

which yields (2.7).

2.2. **Proof of Theorem 1.1.** In order to use the partial Legendre transform, we first recall the definition of modulus of convexity. For a convex function ϕ on \mathbb{R}^n , the modulus of convexity, denoted by m_{ϕ} , is defined by

(2.9)
$$m_{\phi}(t) := \inf\{\phi(x) - \ell_z(x) : |x - z| > t\},\$$

where t > 0 and ℓ_z is a supporting function of ϕ at z. It is clear that m_{ϕ} must be a positive function for a strictly convex function. A result of Heinz [He] implies that in two dimensions, if det $D^2\phi \ge \lambda > 0$, there exists a positive function C(t) > 0depending on λ such that $m_{\phi}(t) \ge C(t) > 0$ (for a more specific C(t), see [Liu, Lemma 2.5]). Now for the partial Legendre transform, we consider the mapping

$$(\xi, \eta) = \mathcal{P}(x_1, x_2) = (\phi_{x_1}, x_2) : B_R(0) \to \mathbb{R}^2.$$

The following important property is revealed in [Liu].

Lemma 2.2 ([Liu, Lemma 2.1]). There exists a constant $\delta > 0$ depending on the modulus of convexity m_{ϕ} defined in (2.9), such that $B_{\delta}(0) \subset \mathcal{P}(B_{R}(0))$.

Remark 2.3. Indeed, from the proof of [Liu, Lemma 2.1], we can see that the dependence on δ only requires the lower bound of m_{ϕ} .

Proof of Theorem 1.1. For any $x \in \Omega$, we denote $R = \frac{\operatorname{dist}(x,\partial\Omega)}{2}$. Without loss of generality, we assume $\mathcal{P}(x) = 0$. Note that ϕ satisfies (1.3), i.e. $\lambda \leq \det D^2 \phi \leq \Lambda$, hence we know that $m_{\phi}(R) \geq C(R) > 0$. By Lemma 2.2, there exists $\delta > 0$ depending on C(R) such that $B_{\delta}(0) \subset \mathcal{P}(B_R(x))$. According to Proposition 2.1, (1.3) and (2.5), we know that \widetilde{u} satisfies (2.7) in $B_{\delta}(0)$ with

$$0 < \lambda \le \det D^2 \phi = -\frac{\phi_{\eta\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \le \Lambda.$$

This means that (2.7) is a uniformly elliptic equation in divergence form.

By the $W^{2,1+\varepsilon}$ -estimate of Monge-Ampère equations [DFS, Sc], there exist $\varepsilon_0 > 0$ depending on λ , Λ , and $C_0 > 0$ depending on R, λ and Λ such that

$$||D^2\phi||_{L^{1+\varepsilon_0}(B_R(x))} \le C_0.$$

Hence, we have

$$\int_{B_{\delta}(0)} (\phi_{\xi\xi}^{\star})^{2+\varepsilon_{0}} d\xi d\eta = \int_{\mathcal{P}^{-1}(B_{\delta}(0))} (\phi_{x_{1}x_{1}})^{-(2+\varepsilon_{0})} \phi_{x_{1}x_{1}} dx_{1} dx_{2}$$

$$= \int_{\mathcal{P}^{-1}(B_{\delta}(0))} (\phi_{x_{1}x_{1}})^{-(1+\varepsilon_{0})} dx_{1} dx_{2}$$

$$= \int_{\mathcal{P}^{-1}(B_{\delta}(0))} \left(\frac{\phi_{x_{2}x_{2}}}{\phi_{x_{1}x_{1}}\phi_{x_{2}x_{2}}}\right)^{1+\varepsilon_{0}} dx_{1} dx_{2}$$

$$\leq \lambda^{-(1+\varepsilon_{0})} \int_{B_{R}(x)} (\phi_{x_{2}x_{2}})^{1+\varepsilon_{0}} dx_{1} dx_{2}$$

$$\leq C\lambda^{-(1+\varepsilon_{0})}.$$

Then by (2.5), we have

$$\int_{B_{\delta}(0)} (-\phi_{\eta\eta}^{\star})^{2+\varepsilon_0} d\xi d\eta \leq \Lambda^{1+\varepsilon_0} \int_{B_{\delta}(0)} (\phi_{\xi\xi}^{\star})^{2+\varepsilon_0} d\xi d\eta \leq C\Lambda^{1+\varepsilon_0} \lambda^{-(1+\varepsilon_0)}.$$

Finally, the standard $W^{2,p}$ theory of uniformly elliptic equations yields

$$\|\phi_{\xi\eta}^{\star}\|_{L^{2+\varepsilon_0}(B_{\delta}(0))} \le C.$$

With the assumptions on F and f, we know that the right-hand sides of (2.7) satisfy

$$\|\widetilde{F}^1 - \widetilde{F}^2 \phi_{\xi\eta}^{\star}\|_{L^{2+\varepsilon_0}(B_{\delta}(0))} \le C, \|\widetilde{F}^2 \phi_{\xi\xi}^{\star}\|_{L^{2+\varepsilon_0}(B_{\delta}(0))} \le C,$$

and

$$\int_{B_{\delta}(0)} \left| \widetilde{f} \phi_{\xi\xi}^{\star} \right|^{\frac{r(2+\varepsilon_{0})}{1+\varepsilon_{0}+r}} d\xi d\eta \leq C \int_{\mathcal{P}^{-1}(B_{\delta}(0))} |f|^{\frac{r(2+\varepsilon_{0})}{1+\varepsilon_{0}+r}} (\phi_{x_{2}x_{2}})^{\frac{(1+\varepsilon_{0})(r-1)}{1+\varepsilon_{0}+r}} dx_{1} dx_{2}
\leq C \left(\int_{B_{R}(x)} |f|^{r} dx_{1} dx_{2} \right)^{\frac{2+\varepsilon_{0}}{1+\varepsilon_{0}+r}} \left(\int_{B_{R}(x)} (\phi_{x_{2}x_{2}})^{1+\varepsilon_{0}} dx_{1} dx_{2} \right)^{\frac{r-1}{1+\varepsilon_{0}+r}}
\leq C \|f\|_{L^{r}(B_{R}(x))}^{\frac{r(2+\varepsilon_{0})}{1+\varepsilon_{0}+r}}.$$

Note that n=2 and $2+\varepsilon_0>2$, $\frac{r(2+\varepsilon_0)}{1+\varepsilon_0+r}>1$ whenever r>1, then the De Giorgi-Nash-Moser's theory [GT, Theorem 8.24] (The original estimate in Theorem 8.24 of [GT] holds in terms of $|\widetilde{u}|_{L^q(B_{\delta}(0))}$ for all q>1. However, by using analogous arguments as in [HL, Page 75], we can extend the validity of this estimate to all q > 0) yields

$$\|\widetilde{u}\|_{C^{\alpha}(B_{\delta/2}(0))} \le C\left(\|\widetilde{u}\|_{L^{\frac{p\varepsilon_0}{1+\varepsilon_0}}(B_{\delta}(0))} + k\right),\,$$

where

$$k = \|\widetilde{F}^{1} - \widetilde{F}^{2} \phi_{\xi \eta}^{\star}\|_{L^{2+\varepsilon_{0}}(B_{\delta}(0))} + \|\widetilde{F}^{2} \phi_{\xi \xi}^{\star}\|_{L^{2+\varepsilon_{0}}(B_{\delta}(0))} + \|\widetilde{f} \phi_{\xi \xi}^{\star}\|_{L^{\frac{r(2+\varepsilon_{0})}{1+\varepsilon_{0}+r}}(B_{\xi}(0))}.$$

Note that by Hölder's inequality there is

$$\left(\int_{B_{\delta}(0)} \widetilde{u}^{\frac{p\varepsilon_{0}}{1+\varepsilon_{0}}} \,\mathrm{d}\xi \,\mathrm{d}\eta\right)^{\frac{1+\varepsilon_{0}}{p\varepsilon_{0}}} = \left(\int_{\mathcal{P}^{-1}(B_{\delta}(0))} u^{\frac{p\varepsilon_{0}}{1+\varepsilon_{0}}} \phi_{x_{1}x_{1}} \mathrm{d}x_{1} \mathrm{d}x_{2}\right)^{\frac{1+\varepsilon_{0}}{p\varepsilon_{0}}} \\
\leq \|u\|_{L^{p}(B_{R}(x))} \cdot \|\phi_{x_{1}x_{1}}\|_{L^{1+\varepsilon_{0}}(B_{R}(x))}^{\frac{1+\varepsilon_{0}}{p\varepsilon_{0}}} \\
\leq C\|u\|_{L^{p}(B_{R}(x))}.$$

Hence, for the original function u, combining with the $C^{1,\alpha}$ estimate of Monge-Ampère equation [Ca] we know that there exists a $\gamma \in (0,1)$ such that

$$||u||_{C^{\gamma}(\mathcal{P}^{-1}(B_{\delta/2}(0)))} \le C\left(||u||_{L^{p}(B_{R}(x))} + ||F||_{L^{\infty}(B_{R}(x))} + ||f||_{L^{r}(B_{R}(x))}\right).$$

By a standard covering argument (see for instance [FR, Remark 2.15]), we know that the estimate is true for any $\Omega' \subset\subset \Omega$, which completes the proof.

2.3. **Proof of Moser-Trudinger type inequality.** In this subsection, we provide the proof of Theorem 1.3.

Moser-Trudinger type inequalities find broad applications in the study of partial differential equations and geometric problems. The classical Moser-Trudinger inequality was initially derived by Trudinger [Tr1], using the power series expansion of the exponential function and Sobolev estimates for individual terms, while carefully examining the dependence on the exponent of the expansions. Subsequently, Moser [Mo2] presented a more direct proof of this inequality and also determined the optimal exponent. Before giving the proof of Theorem 1.3, we first recall the classical Moser-Trudinger inequality.

Theorem 2.4 ([Mo2, Theorem 1]). Let $u \in W_0^{1,n}(\Omega)$ for $n \geq 2$, and

$$\int_{\Omega} |Du|^n \, \mathrm{d}x \le 1.$$

Then there exists a constant C which depends only on n such that

$$\int_{\Omega} e^{\alpha u^p} \mathrm{d}x \le C|\Omega|,$$

where

$$p = \frac{n}{n-1}, \quad \alpha \le \alpha_n := n\omega_{n-1}^{\frac{1}{n-1}},$$

and ω_{n-1} is (n-1)-dimensional surface measure of the unit sphere.

Proof of Theorem 1.3. By the global $W^{2,1+\varepsilon}$ -estimate of Monge-Ampère equations [DFS, Sc, Sa] (with a detailed exposition available in [Le6, Theorem 10.1]), we know that there exist $\varepsilon_0 > 0$ depending on λ and Λ , and $C_0 > 0$ depending only on λ , $\Lambda \|\phi\|_{C^3(\partial\Omega)}$, the uniform convexity radius of $\partial\Omega$ and the C^3 regularity of $\partial\Omega$ such that

$$||D^2\phi||_{L^{1+\varepsilon_0}(\Omega)} \le C_0.$$

Then as the proof of Theorem 1.1, we have

$$\int_{\Omega^*} (\phi_{\xi\xi}^*)^{2+\varepsilon_0} \, \mathrm{d}\xi \mathrm{d}\eta \le C.$$

By (1.3) there is

$$0 < \lambda \le \det D^2 \phi = -\frac{\phi_{\eta\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \le \Lambda.$$

Hence, by (2.4) and (2.8), we obtain

$$||Du||_{\Phi}^{2} = \int_{\Omega} \Phi^{ij} D_{i} u D_{j} u \, dx$$

$$= \int_{\Omega} \phi_{x_{2}x_{2}} u_{x_{1}}^{2} - 2\phi_{x_{1}x_{2}} u_{x_{1}} u_{x_{2}} + \phi_{x_{1}x_{1}} u_{x_{2}}^{2} \, dx_{1} dx_{2}$$

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$$= \int_{\Omega^{\star}} \left(-\frac{\phi_{\eta\eta}^{\star}}{\phi_{\xi\xi}^{\star}} \widetilde{u}_{\xi}^{2} + \widetilde{u}_{\eta}^{2} \right) d\xi d\eta$$

$$\geq \min\{\lambda, 1\} \int_{\Omega^{\star}} \left(\widetilde{u}_{\xi}^{2} + \widetilde{u}_{\eta}^{2} \right) d\xi d\eta$$

$$= \min\{\lambda, 1\} \|D\widetilde{u}\|_{L^{2}(\Omega^{\star})}^{2}.$$

Then there is

$$\int_{\Omega} e^{\beta \frac{u^2}{\|Du\|_{\Phi}^2}} dx \leq \int_{\Omega^{\star}} e^{\beta \min\{\lambda,1\}^{-1} \frac{\tilde{u}^2}{\|D\tilde{u}\|_{L^2(\Omega^{\star})}^2}} \phi_{\xi\xi}^{\star} d\xi d\eta$$

$$\leq \left(\int_{\Omega^{\star}} e^{\beta \min\{\lambda,1\}^{-1} \frac{2+\varepsilon_0}{1+\varepsilon_0} \frac{\tilde{u}^2}{\|D\tilde{u}\|_{L^2(\Omega^{\star})}^2}} d\xi d\eta\right)^{\frac{1+\varepsilon_0}{2+\varepsilon_0}} \left(\int_{\Omega^{\star}} (\phi_{\xi\xi}^{\star})^{2+\varepsilon_0} d\xi d\eta\right)^{\frac{1}{2+\varepsilon_0}}.$$

For any $\beta \leq 4\pi \frac{1+\varepsilon_0}{2+\varepsilon_0} \min\{\lambda,1\}$, by Theorem 2.4 with n=2, we have

$$\int_{\Omega^{\star}} e^{\beta \min\{\lambda,1\}^{-1} \frac{2+\varepsilon_0}{1+\varepsilon_0} \frac{\tilde{u}^2}{\|D\tilde{u}\|_{L^2(\Omega^{\star})}^2}} d\xi d\eta \le C |\Omega^{\star}|.$$

Then by Hölder's inequality,

$$\int_{\Omega} e^{\beta \frac{u^2}{\|Du\|_{\Phi}^2}} dx \leq C |\Omega^{\star}|^{\frac{1+\varepsilon_0}{2+\varepsilon_0}}$$

$$= C \left(\int_{\Omega} \phi_{x_1 x_1} dx_1 dx_2 \right)^{\frac{1+\varepsilon_0}{2+\varepsilon_0}}$$

$$\leq C \left[|\Omega|^{\frac{\varepsilon_0}{1+\varepsilon_0}} \left(\int_{\Omega} (\phi_{x_1 x_1})^{1+\varepsilon_0} dx_1 dx_2 \right)^{\frac{1}{1+\varepsilon_0}} \right]^{\frac{1+\varepsilon_0}{2+\varepsilon_0}}$$

$$\leq C |\Omega|^{\frac{\varepsilon_0}{2+\varepsilon_0}},$$

which completes the proof.

3. Linearized Monge-Ampère equations in higher dimensions

In this section, we establish the proof of Theorem 1.5 in two steps.

3.1. Estimates for linearized Monge-Ampère equations. In this subsection, we prove a weak maximum principle for linearized Monge-Ampère equations, which will be used in the proof of Theorem 1.5.

Theorem 3.1 (Weak maximum principle). Let $\phi \in C^2(\Omega)$ be a convex function satisfying (1.3). For q > n, let $F : \Omega \to \mathbb{R}^n$ be a vector field satisfying $||F_{\phi}||_{L^q(\Omega)} < \infty$, where $F_{\phi}(x) := (D^2\phi(x))^{1/2} F$ and $f : \Omega \to \mathbb{R}$ be a function satisfying $||f||_{L^{q_*}(\Omega)} < \infty$, where $q_* = \frac{nq}{n+q}$. For every solution u to

(3.1)
$$D_{i}\left(\Phi^{ij}D_{i}u\right) \geq \operatorname{div}F + f \quad in \quad \Omega,$$

we have

$$\sup_{\Omega} u \leq \sup_{\partial \Omega} u^{+} + C\left(n, q, \lambda, \Lambda, \operatorname{diam}(\Omega)\right) \left(\|F_{\phi}\|_{L^{q}(\Omega)} + \|f\|_{L^{q_{*}}(\Omega)} \right) |\Omega|^{\frac{1}{n} - \frac{1}{q}},$$

where $u^+ := \max\{u, 0\}$.

As a corollary, we have

Corollary 3.2 (Global estimate for solutions to the Dirichlet problem). Let $\phi \in C^2(\Omega)$ be a convex function satisfying (1.3). For q > n, let $F : \Omega \to \mathbb{R}^n$ be a vector field satisfying $||F_{\phi}||_{L^q(\Omega)} < \infty$, where $F_{\phi}(x) := (D^2\phi(x))^{1/2}F$ and $f : \Omega \to \mathbb{R}$ be a function satisfying $||f||_{L^{q_*}(\Omega)} < \infty$, where $q_* = \frac{nq}{n+q}$. For every section $S_{\phi}(x_0, h)$ with $S_{\phi}(x_0, h_0) \subset\subset \Omega$ for $h_0 \geq h$ and every solution u to

$$\begin{cases} D_j \left(\Phi^{ij} D_i u \right) = \operatorname{div} F + f & in \quad S_{\phi}(x_0, h), \\ u = 0 & on \quad \partial S_{\phi}(x_0, h), \end{cases}$$

we have

$$\sup_{S_{\phi}(x_0,h)} |u| \le C(n,q,\lambda,\Lambda,\operatorname{diam}(\Omega),h_0) \left(||F_{\phi}||_{L^q(S_{\phi}(x_0,h))} + ||f||_{L^{q_*}(S_{\phi}(x_0,h))} \right) h^{\frac{1}{2} - \frac{n}{2q}}.$$

We will use De Giorgi's iteration to prove Theorem 3.1 and Corollary 3.2, and De Giorgi's iteration is a very powerful tool for dealing with elliptic equations in divergence form. It is usually reduced to the following iteration lemma, the proof of which can be found in [CW].

Lemma 3.3 ([CW, Lemma 4.1]). Let $\omega(t)$ be a nonnegative and nonincreasing function in an interval $[k_0, +\infty)$. Suppose that there holds for all $h > k \ge k_0$,

$$\omega(h) \le \frac{C}{(h-k)^{\alpha}} [\omega(k)]^{\beta},$$

where $\alpha > 0$ and $\beta > 1$. Then we have

$$\omega(k_0+d)=0,$$

where

$$d = C^{\frac{1}{\alpha}} [\omega(k_0)]^{\frac{\beta-1}{\alpha}} 2^{\frac{\beta}{\beta-1}}.$$

First, we give the proof of Theorem 3.1, and it follows the same idea as the proof of [CW, Theorem 4.2].

Proof of Theorem 3.1. Denote $l = \sup_{\partial\Omega} u^+$. Consider $v = (u - k)^+$ for $k \ge l$. Note that v = u - k, Dv = Du a.e. in $\{u > k\}$ and v = 0, Dv = 0 a.e. in $\{u \le k\}$. Taking v as test function in (3.1), we have

$$-\int_{\Omega} \Phi^{ij} D_i u D_j v \, dx \ge -\int_{\Omega} F_i D_i v \, dx + \int_{\Omega} f v \, dx.$$

Then

$$\int_{\Omega} \Phi^{ij} D_{i} v D_{j} v \, dx \leq \int_{\Omega} F_{i} D_{i} v \, dx - \int_{\Omega} f v \, dx
= \int_{\Omega} \left(D^{2} \phi \right)^{1/2} F \cdot \left(D^{2} \phi \right)^{-1/2} D v \, dx - \int_{\Omega} f v \, dx
\leq \left(\|F_{\phi}\|_{L^{q}(\Omega)} \|D^{\phi} v\|_{L^{2}(\Omega)} + \|f\|_{L^{q_{*}}(\Omega)} \|v\|_{L^{2^{*}}(\Omega)} \right) |A(k)|^{\frac{1}{2} - \frac{1}{q}},$$

where $q_* = \frac{nq}{n+q}$, $2^* = \frac{2n}{n-2}$ (we only consider $n \geq 3$ here, and n=2 is similar),

$$D^{\phi}v := (D^2\phi)^{-1/2} Dv$$
 and $A(k) = \{x \in \Omega \mid u(x) > k\}.$

Note that by (1.3) and $\Phi = (\det D^2 \phi)(D^2 \phi)^{-1}$, there is

$$||D^{\phi}v||_{L^{2}(\Omega)} = \left(\int_{\Omega} \left[(D^{2}\phi)^{-1} \right]^{ij} D_{i}v D_{j}v \, dx \right)^{\frac{1}{2}} \le \lambda^{-1/2} \left(\int_{\Omega} \Phi^{ij} D_{i}v D_{j}v \, dx \right)^{\frac{1}{2}}.$$

Denote

$$F_0 = \lambda^{-1/2} ||F_{\phi}||_{L^q(\Omega)} + C_{Sob} ||f||_{L^{q_*}(\Omega)},$$

where C_{Sob} is the Sobolev constant in Theorem 1.2. Hence, we know by Cauchy-Schwarz's inequality and Theorem 1.2 that

$$\int_{\Omega} \Phi^{ij} D_{i} v D_{j} v \, dx \leq \left(\|F_{\phi}\|_{L^{q}(\Omega)} \|D^{\phi} v\|_{L^{2}(\Omega)} + \|f\|_{L^{q_{*}}(\Omega)} \|v\|_{L^{2^{*}}(\Omega)} \right) |A(k)|^{\frac{1}{2} - \frac{1}{q}}$$

$$\leq F_{0} \left(\int_{\Omega} \Phi^{ij} D_{i} v D_{j} v \, dx \right)^{\frac{1}{2}} |A(k)|^{\frac{1}{2} - \frac{1}{q}}$$

$$\leq \frac{1}{2} \int_{\Omega} \Phi^{ij} D_{i} v D_{j} v \, dx + \frac{F_{0}^{2}}{2} |A(k)|^{1 - \frac{2}{q}},$$

which implies

$$\int_{\Omega} \Phi^{ij} D_i v D_j v \, dx \le F_0^2 |A(k)|^{1 - \frac{2}{q}},$$

i.e.

$$\left(\int_{\Omega} \Phi^{ij} D_i v D_j v \, dx \right)^{\frac{1}{2}} \le F_0 |A(k)|^{\frac{1}{2} - \frac{1}{q}}.$$

By Theorem 1.2, there is

$$||v||_{L^{2^*}(\Omega)} \le C_{Sob} F_0 |A(k)|^{\frac{1}{2} - \frac{1}{q}},$$

Note that $v = (u - k)^+$. Thus, when h > k, there is

$$||v||_{L^{2^*}(\Omega)} \ge (h-k)|A(h)|^{\frac{1}{2^*}}.$$

Hence, we have

$$|A(h)| \le \frac{(C_{Sob}F_0)^{2^*}}{(h-k)^{2^*}} |A(k)|^{\frac{n(q-2)}{q(n-2)}}, \quad \forall h > k \ge l.$$

Applying Lemma 3.3 with $\alpha=2^*$ and $\beta=\frac{n(q-2)}{q(n-2)}>1$ yields

$$A(l+d) = 0,$$

where

$$d = C_{Sob}F_0|A(0)|^{\frac{1}{n} - \frac{1}{q}} \cdot 2^{\frac{n(q-2)}{2(q-n)}}.$$

Therefore, we obtain that

$$\sup_{\Omega} u \le l + d \le \sup_{\partial \Omega} u^{+} + C \left(\|F_{\phi}\|_{L^{q}(\Omega)} + \|f\|_{L^{q_{*}}(\Omega)} \right) |\Omega|^{\frac{1}{n} - \frac{1}{q}}.$$

Next, we prove Corollary 3.2.

Proof of Corollary 3.2. Using (1.3) and the volume estimates for sections (see, for example, [Fi, Lemma 4.6]), we obtain:

$$c_1(n,\lambda,\Lambda)h^{n/2} \le |S_{\phi}(x_0,h)| \le C_1(n,\lambda,\Lambda)h^{n/2}$$

Hence, applying Theorem 3.1 with $\Omega = S_{\phi}(x_0, h)$ and u = 0 on $\partial S_{\phi}(x_0, h)$, we have

$$\sup_{S_{\phi}(x_0,h)} u \le C \left(\|F_{\phi}\|_{L^q(S_{\phi}(x_0,h))} + \|f\|_{L^{q_*}(S_{\phi}(x_0,h))} \right) h^{\frac{1}{2} - \frac{n}{2q}}.$$

For -u we can similarly show

$$\sup_{S_{\phi}(x_0,h)} (-u) \le C \left(\|F_{\phi}\|_{L^q(S_{\phi}(x_0,h))} + \|f\|_{L^{q_*}(S_{\phi}(x_0,h))} \right) h^{\frac{1}{2} - \frac{n}{2q}}.$$

Combining the two inequalities gives

$$\sup_{S_{\phi}(x_0,h)} |u| \le C(n,q,\lambda,\Lambda,\operatorname{diam}(\Omega),h_0) \left(||F_{\phi}||_{L^q(S_{\phi}(x_0,h))} + ||f||_{L^{q_*}(S_{\phi}(x_0,h))} \right) h^{\frac{1}{2} - \frac{n}{2q}}.$$

3.2. **Proof of Theorem 1.5.** The proof of Theorem 1.5 is followed by the combination of Corollary 3.2 with Caffarelli-Gutiérrez's Harnack inequality. For reader's convenient, we recall the Harnack inequality here.

Theorem 3.4 ([CG, Theorem 5]). Assume $n \geq 2$. Let $\phi \in C^2(\Omega)$ be a convex function satisfying (1.3). Let $u \in W^{2,n}_{loc}(\Omega)$ be a nonnegative solution of the homogeneous linearized Monge-Ampère equation

$$\Phi^{ij}D_{ij}u = 0$$

in a section $S_{\phi}(x_0, 2h) \subset\subset \Omega$. Then there is

$$\sup_{S_{\phi}(x_0,h)} u \le C(n,\lambda,\Lambda) \inf_{S_{\phi}(x_0,h)} u.$$

The proof of Theorem 1.5 mirrors that of Theorem 1.3 in [Le2, P284-P285]. For the sake of completeness, we outline it briefly here.

Sketch of the proof of Theorem 1.5. By noticing the $C^{1,\alpha}$ of standard Monge-Ampère equations provided (1.3), we know it suffices to show for $h \leq h_0$, there is

$$(3.2) \qquad \underset{S_{\phi}(x_0,h)}{\operatorname{osc}} u \leq C \left(\|u\|_{L^{\infty}(S_{\phi}(x_0,h))} + \|F_{\phi}\|_{L^{q}(S_{\phi}(x_0,h))} + \|f\|_{L^{q_*}(S_{\phi}(x_0,h))} \right) h^{\gamma_0},$$

where $\gamma_0 \in (0,1)$ depending only on n, q, λ and Λ . To prove (3.2), we split u = v + w where

$$\begin{cases} \Phi^{ij} D_{ij} v = \operatorname{div} F + f & \text{in } S_{\phi}(x_0, h), \\ v = 0 & \text{on } \partial S_{\phi}(x_0, h), \end{cases}$$

and

$$\begin{cases} \Phi^{ij} D_{ij} w = 0 & \text{in } S_{\phi}(x_0, h), \\ w = u & \text{on } \partial S_{\phi}(x_0, h). \end{cases}$$

Then, by Corollary 3.2,

$$\sup_{S_{\phi}(x_0,h)} |v| \le C \left(\|F_{\phi}\|_{L^q(S_{\phi}(x_0,h))} + \|f\|_{L^{q_*}(S_{\phi}(x_0,h))} \right) h^{\frac{1}{2} - \frac{n}{2q}}.$$

By Caffarelli-Gutiérrez's Harnack inequality (Theorem 3.4),

$$\underset{S_{\phi}(x_0,h/2)}{\operatorname{osc}} w \le \beta \underset{S_{\phi}(x_0,h)}{\operatorname{osc}} w.$$

Therefore, we have

$$\begin{aligned} & \underset{S_{\phi}(x_0,h/2)}{\operatorname{osc}} \, u \leq \underset{S_{\phi}(x_0,h/2)}{\operatorname{osc}} \, w + \underset{S_{\phi}(x_0,h/2)}{\operatorname{osc}} \, v \\ & \leq \beta \, \underset{S_{\phi}(x_0,h)}{\operatorname{osc}} \, w + 2 \| v \|_{L^{\infty}(S_{\phi}(x_0,h))} \\ & \leq \beta \, \underset{S_{\phi}(x_0,h)}{\operatorname{osc}} \, u + C \left(\| F_{\phi} \|_{L^q(S_{\phi}(x_0,h))} + \| f \|_{L^{q_*}(S_{\phi}(x_0,h))} \right) h^{\frac{1}{2} - \frac{n}{2q}}. \end{aligned}$$

Finally, by a standard iteration, there is

$$\underset{S_{\phi}(x_0,h)}{\operatorname{osc}} u \leq C \left(\frac{h}{h_0} \right)^{\gamma_0} \left(\|u\|_{L^{\infty}(S_{\phi}(x_0,h))} + \left(\|F_{\phi}\|_{L^q(S_{\phi}(x_0,h))} + \|f\|_{L^{q_*}(S_{\phi}(x_0,h))} \right) h_0^{\frac{1}{2} - \frac{n}{2q}} \right)$$

for some $\gamma_0 \in (0,1)$ and some constant C > 0 depending only on $n, q, \lambda, \Lambda, h_0$ and $\operatorname{diam}(\Omega)$, which yields (3.2).

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